

#### LA-UR-17-28210

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Title: Advancing a phase field dislocation dynamics code to model HCP

materials

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Intended for: Lightening Talk (Student Presentation) Wed. 9/13/17 at 10am in CNLS

room 102 Report

Issued: 2017-09-12



# Advancing a phase field dislocation dynamics code to model HCP materials

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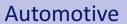
<sup>&</sup>lt;sup>3</sup> XCP-1: Lagrangian Codes

<sup>&</sup>lt;sup>4</sup> T-3: Fluid Dynamics and Solid Mechanics

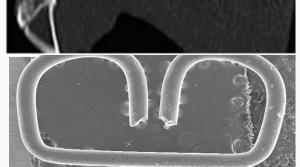
<sup>&</sup>lt;sup>5</sup> MST-8: Materials Science in Radiation & Dynamics Extremes

## Motivation: HCP metals and their alloys

 High strength to weight ratio, biocompatibility, radiation resistance, fatigue resistance



## Biomedical







# Phase Field Dislocation Dynamics PFDD Model

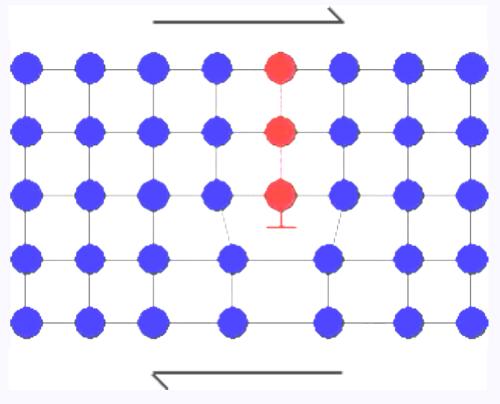
Phase field frameworks describe physical behavior by tracking one or more scalar order parameters(ζ) and evolving them through the minimization of the system's total energy (E), equilibrating between every time step.

$$\frac{\delta E(\zeta)}{\delta \zeta} = 0$$

$$E = E^{strain} + E^{core}$$

• When a dislocation slips the atoms change neighbors and the order parameter takes on a value of 1. The order parameter remains 0 if no slip occurs.

## **Perfect Edge Dislocation Movement**



www.youtube.com/watch?v=iKKxTP6xp74

$$\epsilon_{ij}^{p}(\mathbf{x},t) = \frac{1}{2} \sum_{\alpha=1}^{N} b\zeta_{\alpha}(\mathbf{x},t) \delta_{n}(s_{i}^{\alpha} m_{j}^{\alpha} + s_{j}^{\alpha} m_{i}^{\alpha})$$

- Plasticity or deformation is mediated by the motion and interaction of dislocations
- Thus the **plastic strain**,  $\epsilon_{ij}^p(x,t)$ , will be directly proportional to the number of gliding dislocations which are our phase field variable,  $\zeta$  (order parameter).

N = number of slip systems

 $\alpha$  = a specific slip system

 $s^{\alpha}$  = slip direction

 $m^{\alpha}$  = normal to the slip plane

 $\delta_n$  = Dirac distribution supported on the slip plane, n

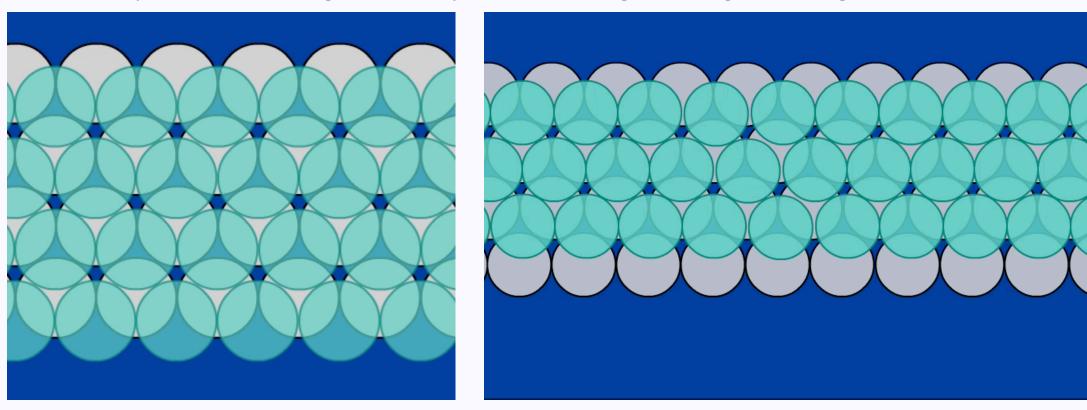
b = the magnitude of the Burgers vector

$$\hat{A}_{mnuv}(\mathbf{k}) = C_{mnuv} - C_{kluv}C_{ijmn}\hat{G}_{ki}(k)k_jk_l$$

$$E^{strain}(\zeta) = \frac{1}{2} \int \hat{A}_{mnuv}(\mathbf{k}) \hat{\epsilon}_{mn}^{p}(\mathbf{k}) \hat{\epsilon}_{mn}^{p*}(\mathbf{k}) \frac{d^{3}k}{(2\pi)^{3}} - \int \sigma_{ij}^{appl} \epsilon_{ij}^{p} d^{3}x$$

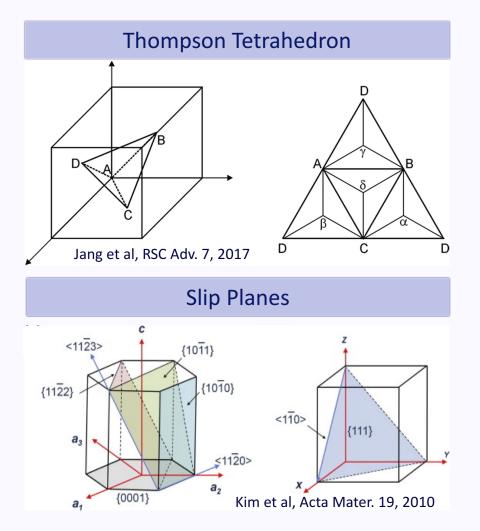
# Partial dislocations & stacking fault widths

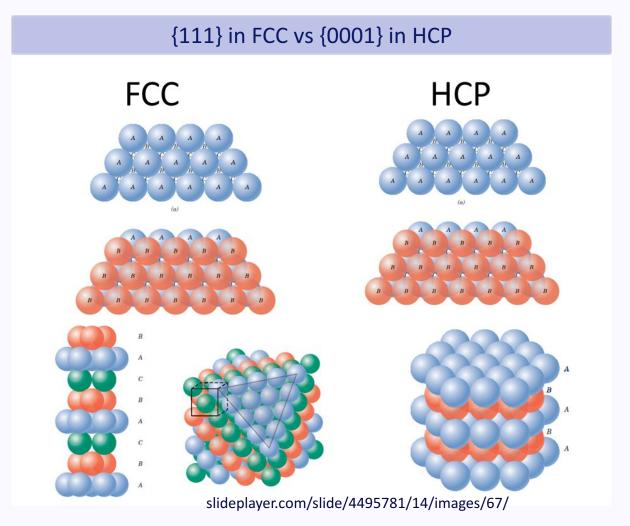
It is energetically favorable for a perfect dislocation to split into two **partial dislocations**. These two partials are like signed and repel each other, generating a **stacking fault width**.



www.princeton.edu/~maelabs/mae324/07/07mae\_52a.htm

# Face Centered Cubic & Hexagonal Close Packed

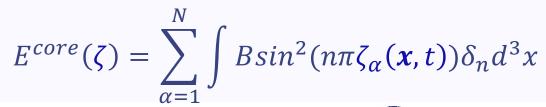


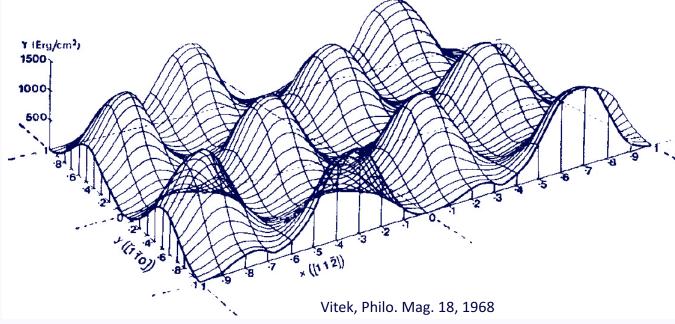


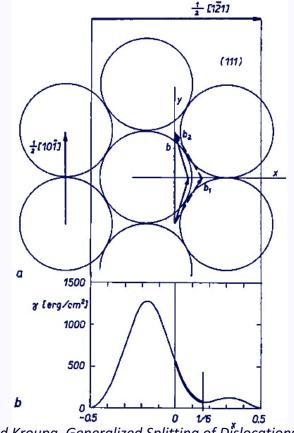
# Core energy and y-surface

## **Perfect Dislocations**

## **Partial Dislocations**

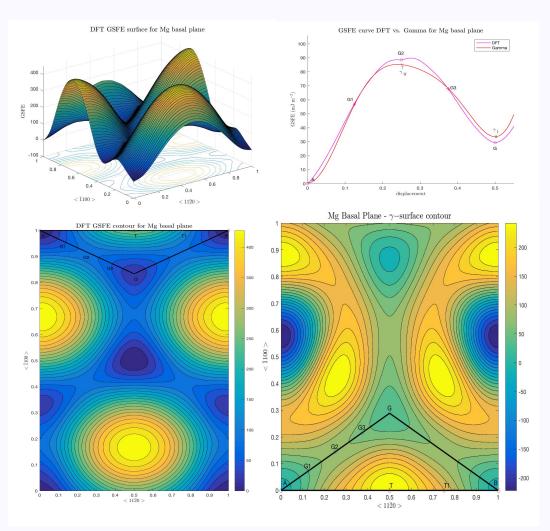






Vitek and Kroupa, Generalized Splitting of Dislocations, 1968

## Magnesium basal plane {0001} and γ-surface



 $\Delta$ AGB: vectors for a perfect and 2 partial dislocation **Schoeck** parameterization uses 7 GSFE values taken from the  $\Delta$ AGB (DFT) to generate a  $\gamma$ -surface that can be input into the PFDD code.

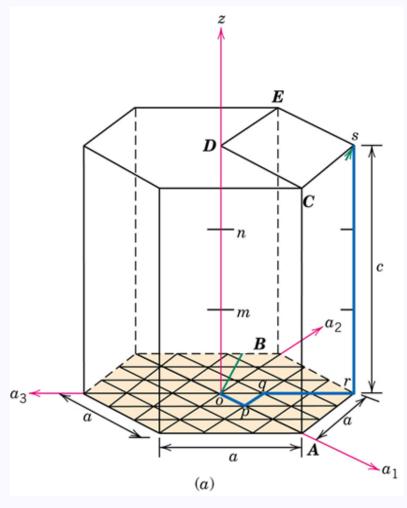
```
 \gamma[x,y] = c1*(\cos(2p^*y) + \cos(p^*y + q^*x) + \cos(p^*y - q^*x)) 
 + c2*(\cos(2q^*x) + \cos(3p^*y + q^*x) + \cos(3p^*y - q^*x)) 
 + c3*(\cos(4p^*y) + \cos(2p^*y + 2q^*x) + \cos(-2p^*y + 2q^*x)) 
 + c4*(\cos(p^*y + 3q^*x) + \cos(-p^*y + 3q^*x) + \cos(4p^*y + 2q^*x) 
 + \cos(-4p^*y + 2q^*x) + \cos(5p^*y + q^*x) + \cos(5p^*y - q^*x)) 
 + a1*(\sin(2p^*y) - \sin(p^*y + q^*x) + \sin(-p^*y + q^*x)) 
 + a3*(\sin(4p^*y) - \sin(2p^*y + 2q^*x) + \sin(-2p^*y + 2q^*x))
```

#### Constants:

```
c0 = 0.823*(4*G-6*G1+6*G2-7.392*G3+0.804*T+0.804*T1) \\ c1 = 0.274*(-8*G+12*G1-12*G2+14.785*G3-1.608*T+0.215*T1) \\ c2 = 0.091(23.072*G-29.138*G1+32.785*G2-42.215*G3+2.569*T-2.412*T1) \\ c3 = 0.137*(-8*G+12*G1-12*G2+14.785*G3+0.215*T-1.608*T1) \\ c4 = 0.023*(1.856*G-13.723*G1+6.431*G2-4.277*G3-0.962*T+3.531*T1) \\ a1 = 0.137*(-32*G+48*G1-48*G2+62.785*G3-4.608*T-2.785*T1) \\ a3 = 0.046*(17.072*G-19.292*G1+31.923*G2-34.708*G3+3.341*T-8.354*T1) \\ \end{array}
```

Schoeck, Philo. Mag. A 81, 2009

## HCP: Indices conversion & Lattice Rotation



## **Indices Conversions**

- Miller-Bravais [uvtw]
- Miller [*UVW*]

$$U = u - t$$

$$V = v - t$$

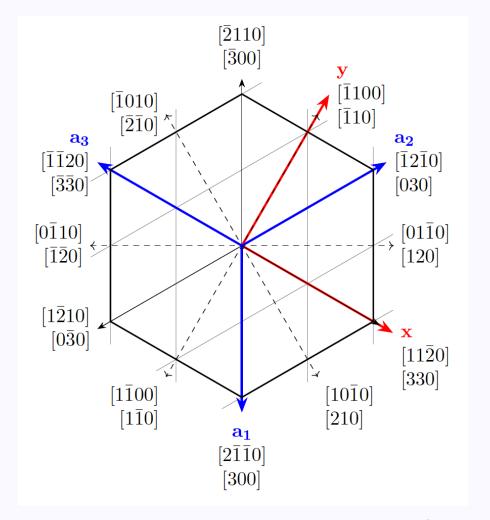
$$W = w$$

Normalize by magnitude

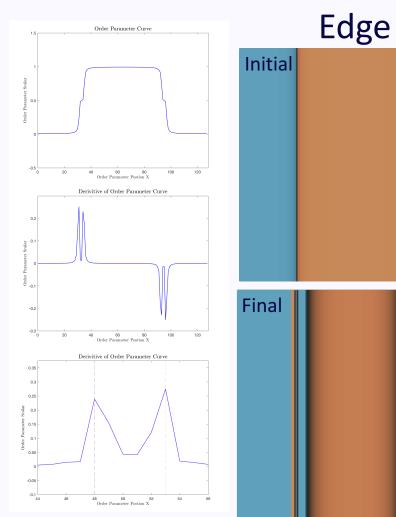
## **Lattice Rotation**

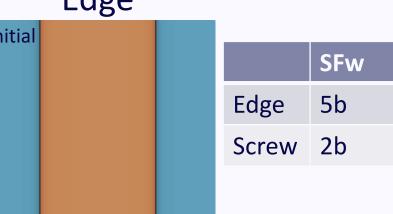
- x-axis:  $[11\bar{2}0] \rightarrow [100]$
- y-axis:  $[\bar{1}100] \to [010]$

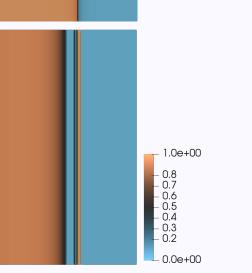
$$[R_B]_I = \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & 0\\ \frac{-\sqrt{3}}{2} & \frac{\sqrt{3}}{2} & 0\\ 0 & 0 & 1 \end{bmatrix}$$

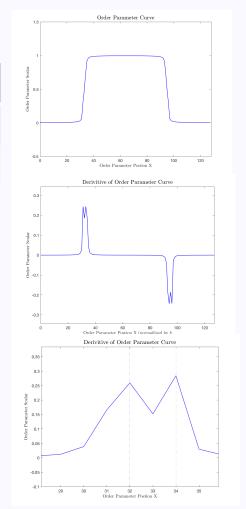


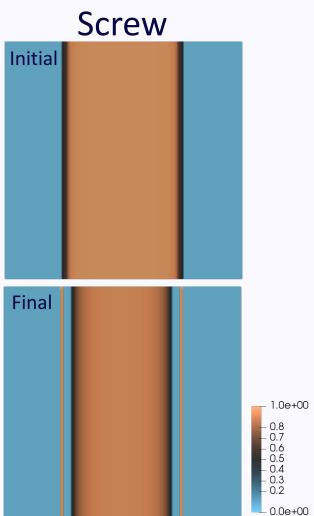
## Results







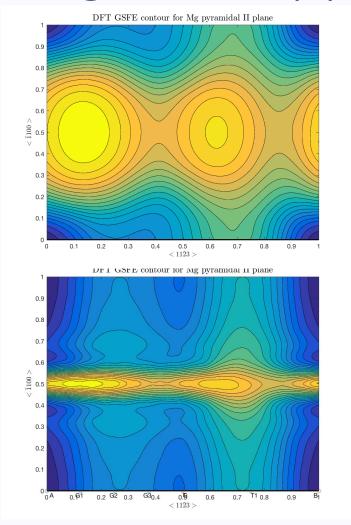


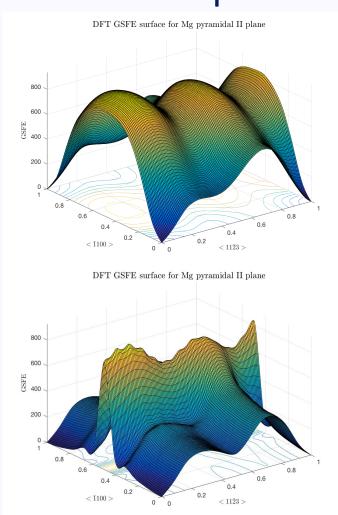


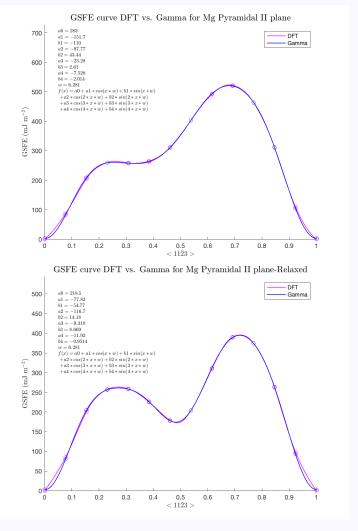
# Comparison of calculations for SF width

Mg Basal Stacking Fault Widths					
Author	EDGE	SCREW	b (Å)	Method	Additional
Weaver	5 b	2 b	3.19	PFDD	Schoeck Parameterization
Yasi et al	5.2 b	2.0 b	3.2	Ab initio	VASP, LGF, GGA with P&W exchange-correlation potential flexible B.C.
	4.5 b	2.0 b	3.2	Sun EAM	periodic in dislocation line direction
	4.0 b	0.4 b	3.2	Liu EAM	periodic in dislocation line direction
Yin et al	7.5 b	3.5 b	3.186	AniLinElastTheory	Lattice/ elastic constants from experiements, SFE from DFT
	7.0 b	4.0 b	3.189	Ab initio	VASP, GGA with PBE parameterization
Wu et al	2.20 b	1.26 b	3.187	Ab initio	DFT
	3.92 b	1.26 b	3.187	MEAM	
Fan et al	8.5 b	4.38 b	3.2	Peierls-Nabarro	fits surface with reciprocal lattice vectors to a 2D Fourier series
Shen et al	5.88 b	2.16 b	*3.2	EAM	
Groh et al	8.0 b	5.13 b	*3.2	EAM	
Wang et al	6.66 b		*3.2	Peierls-Nabarro	

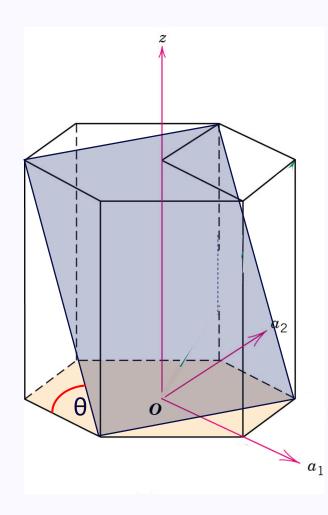
# Magnesium pyramidal II plane γ-surface







# HCP: Lattice rotation for Pyramidal II



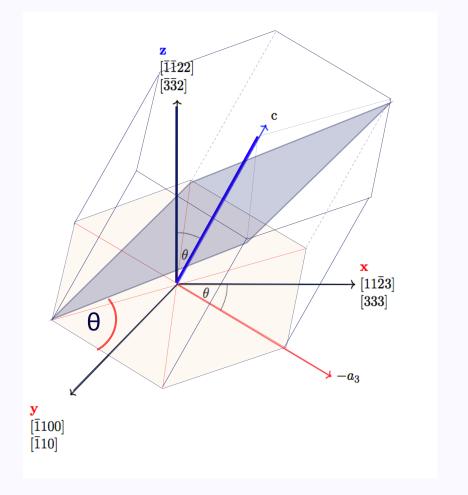
### **Lattice Rotation**

- x-axis:  $[11\bar{2}3] \rightarrow [100]$
- y-axis:  $[\bar{1}100] \rightarrow [010]$
- We can apply an additional rotation to the basal plane rotation,  $[R_B]_I$ , around the y axis by  $\theta$  to get:

$$[R_P]_I = \begin{bmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{bmatrix} [R_B]_I$$

$$\cos \theta = \frac{a^2}{\sqrt{a^2 + c^2}} \qquad \sin \theta = \frac{c^2}{\sqrt{a^2 + c^2}}$$

$$[R_P]_I = \begin{bmatrix} \frac{a}{2\sqrt{a^2 + c^2}} & \frac{a}{2\sqrt{a^2 + c^2}} & \frac{c}{\sqrt{a^2 + c^2}} \\ \frac{-\sqrt{3}}{2} & \frac{\sqrt{3}}{2} & 0 \\ -\frac{c}{2\sqrt{a^2 + c^2}} & -\frac{c}{2\sqrt{a^2 + c^2}} & \frac{a}{\sqrt{a^2 + c^2}} \end{bmatrix}$$



## Whats Next?

- Pyramidal II:
  - Add 1D fit equation to the code for preliminary results
  - Fit for the entire 2D surface and add to code
- Take a look at the Prismatic slip planes
  - Conduct a literature review of dislocation dissociation on this plane
  - Fit the gamma surface
  - Create a rotation matrix for easy input into code
- Implement for other HCP materials and Mg alloys
  - Generalize the gamma surface parameterization for each slip plane so only a few GSFE points from the surface at key points are required for the code.
- Begin to look at other microstructural influences (i.e. twinning, dislocation nucleation, etc.) on plastic deformation within HCP materials

## Material constants and Equations

#### Used

```
a = 3.19E-10;
b = 3.19E-10;
mu = 18.0E9;
young = 46.8E9;
c0 = 112.6205E-3;
c1 = 0.7200E-3;
c2 = -58.3781E-3;
c3 = 28.1313E-3;
c4 = -3.9212E-3;
a1 = 25.2325E-3;
a3 = -23.8545E-3;
isf = 29.3216E-3;
usf = 88.4906E-3;
```

#### **DTF** Derived

```
C12 = 27E9;
C11 = 63E9;
C44 = 0.5(C11-C12) = 18E9;
|| = C12 = 27E9;
mu = C44 = 0.5(C11-C12) = 18E9;
voung =
mu(3.0*II+2.0*mu)/(mu+II) =
46.8E9;
nu = young/2.0/(mu)-1.0 = 0.3;
S11 = 1.0/(young) = 2.14E-11;
S12 = -nu/(young) = -6.41E-12;
```

S44 = 2.0\*(S11-S12) = 5.56E-11;

#### Simmons&Wang

C12 = 25.94E9;

```
C11 = 63.48E9;
C44 = 0.5*(C11-C12) = 18.77E9;
II = C12 = 25.94E9;
mu = C44 = 18.77E9;
young = mu(3.0*II+2.0*mu)/(mu+II) =
48.43E9;
nu = young/2.0/(mu)-1.0 =
0.2901;
S11 = 1.0/(young) = 2.065E-11;
S12 = -nu/(young) = -5.990E-12;
S44 = 2.0(S11-S12) = 5.328E-11;
```

#### **Equations**

```
C44 = mu;

nu = young/2.0/(mu)-1.0;

C12 = 2.0*nu*C44/(1.0-2.0*nu);

C11 = 2.0*C44+C12;

II = C12;

S11 = 1.0/(young);

S12 = -nu/(young);

S44 = 2*(S11-S12);
```